

# Hadron Yield Correlation in Quark Combination Models in High-Energy AA Collisions

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We study the hadron yield correlation in the combination models in high-energy heavy-ion collisions. We derive the relationship between the average yields of different hadrons produced in the combination of a system consisting of equal numbers of quarks and antiquarks. We present the results for the directly produced hadrons as well as those for the final hadrons including the strong and electromagnetic decay contributions. We also study the net quark influence by considering the case when the number of quarks is larger than that of antiquarks. We make comparison with the data wherever possible.

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## I. INTRODUCTION

Hadron yield correlations, measured by the ratios of the average yields of different hadrons produced in high-energy reactions, are one kind of characteristic properties of hadronization mechanisms. It is usually expected that these correlations are more or less independent of the particular model especially when the combination mechanism is concerned. Such properties were therefore considered as a good probe for the hadronization mechanism in different high-energy reactions already in the 1970s to 1990s [1–8]. The study of these correlations has attracted much attention in heavy-ion collisions recently [9–15] because they are considered as a probe to test whether the (re-)combination mechanism is at work. This is interesting because whether the combination mechanism works might be considered as one of the signatures for the formation of the bulk color-deconfined quark matter system before the hadronization takes place. Experimental results are available from the Relativistic Heavy Ion Collider (RHIC) [16, 17], from relatively low-energy collisions such as those obtained by NA49, NA61, and CBM Collaborations at the Super Proton Synchrotron (SPS) [17–20], and more recently from the very high energy reactions at the Large Hadron Collider (LHC) [21, 22]. These results seem to suggest a dramatic change for collisions from lower to higher energies, which is considered as one of the hints for phase transition [18].

In order to make a judgment whether the combination mechanism is at work by comparing the theoretical results with the corresponding experimental data, it is important to see whether, if so, to what extent, the theoretical results depend on the particular model used in obtaining these results. There are many studies that have been made in the literature [9–13, 23]. These studies are usually based on some particular (re-)combination or coalescence models and/or some particular assumptions. It is not clear whether the results obtained depend on the particular assumption(s) made in these particular models. For this purpose, in this paper, we will make a systematic study of the average yields of different identified hadrons and their relationships obtained in the combination mechanism. We will start the study by considering the case for the combi-

nation of a system of quarks and antiquarks from the basic ideas of the combination mechanism. We will make the study as independent of the particular models as possible but present the assumptions and/or inputs explicitly whenever necessary.

The rest of the paper is organized as follows. In Sec. II, we derive the formulas for calculating the average yields of hadrons and their relationships in the combination of a system of quarks and antiquarks. We consider a system where the number of quarks and that of antiquarks are equal and discuss the net quark influence as well. We compare the results with the available data in Sec. III. Here, the experimental results from the LHC [21, 22] are taken as an example to test the predictions for the case where the net quark influence is considered as negligible, while those from the RHIC and the SPS [18, 24–36] are used to test the net quark influence. A short summary is presented in Sec. V.

## II. HADRON YIELD RATIOS IN THE COMBINATION MODELS

In this section, we begin with the general formalism of hadron yields in the combination models based on the basic ideas. For this purpose, we start with a quark-antiquark system as general as possible. Then we simplify the results by using some explicit assumptions, simplifications, and/or approximations. We present the results for the ratios of the yields of hadrons directly produced as well as those including the contributions from the resonance decays.

### A. The general formalism

We start with the most general case and consider a system of  $N_q$  quarks and  $N_{\bar{q}}$  antiquarks. We denote the number of quarks of flavor  $q_i$  by  $N_{q_i}$  so that  $\sum_i N_{q_i} = N_q$  and similarly  $\sum_i N_{\bar{q}_i} = N_{\bar{q}}$ . These quarks and antiquarks combine with each other to form the color singlet hadrons. The number density of the directly produced hadrons is given

by

$$f_{M_j}(p_{M_j}) = \sum_{q_1 \bar{q}_2} \int dp_1 dp_2 f_{q_1 \bar{q}_2}(p_1, p_2; N_{q_i}, N_{\bar{q}_i}) \times \mathcal{R}_{M_j, q_1 \bar{q}_2}(p_{M_j}, p_1, p_2; N_{q_i}, N_{\bar{q}_i}), \quad (1)$$

$$f_{B_j}(p_{B_j}) = \sum_{q_1 q_2 q_3} \int dp_1 dp_2 dp_3 \times f_{q_1 q_2 q_3}(p_1, p_2, p_3; N_{q_i}, N_{\bar{q}_i}) \times \mathcal{R}_{B_j, q_1 q_2 q_3}(p_{B_j}, p_1, p_2, p_3; N_{q_i}, N_{\bar{q}_i}), \quad (2)$$

where  $f_{M_j}$  and  $f_{B_j}$  are the momentum distributions for the produced meson  $M_j$  and baryon  $B_j$ , respectively;  $f_{q_1 \bar{q}_2}(p_1, p_2; N_{q_i}, N_{\bar{q}_i})$  and  $f_{q_1 q_2 q_3}(p_1, p_2, p_3; N_{q_i}, N_{\bar{q}_i})$  are the two- and three-particle joint momentum distributions for  $(q_1 \bar{q}_2)$  and  $(q_1 q_2 q_3)$ , respectively. The kernel functions  $\mathcal{R}_{M_j, q_1 \bar{q}_2}(p_{M_j}, p_1, p_2; N_{q_i}, N_{\bar{q}_i})$  and  $\mathcal{R}_{B_j, q_1 q_2 q_3}(p_{B_j}, p_1, p_2, p_3; N_{q_i}, N_{\bar{q}_i})$  stand for the probability density for  $q_1$  and  $\bar{q}_2$  with momenta  $p_1$  and  $p_2$  to combine into a meson  $M_j$  of momentum  $p_{M_j}$  and that for  $q_1, q_2$  and  $q_3$  with momenta  $p_1, p_2$ , and  $p_3$  to coalesce into a baryon  $B_j$  of momentum  $p_{B_j}$ . Here, in the arguments, we use  $N_{q_i}$  and  $N_{\bar{q}_i}$  to represent the dependence of these functions on the numbers of the quarks and antiquarks of different flavors, and also on the total collision energy  $\sqrt{s}$  of AA reactions. We note in particular that not only the joint distributions  $f_{q_1 \bar{q}_2}$  and  $f_{q_1 q_2 q_3}$  but also the probability densities  $\mathcal{R}_{M_j, q_1 \bar{q}_2}$  and  $\mathcal{R}_{B_j, q_1 q_2 q_3}$  are in general dependent on  $N_{q_i}$  and  $N_{\bar{q}_i}$ . This is because, for finite  $N_{q_i}$  and  $N_{\bar{q}_i}$ , the probability for a given quark  $q_1$  to combine with a specified antiquark  $\bar{q}_2$  to form a specified meson  $M_j$  or two other specified quarks  $q_2 q_3$  to form a baryon  $B_j$  is in general dependent on the number  $N_{q_i}$  of existing quarks of different flavors and the number  $N_{\bar{q}_i}$  of antiquarks.

We note in particular the relationship between the description presented here and those given in the literature in different models based on the combination mechanism such as the coalescence model [9, 10], the recombination model [11, 12], and the quark combination model [13, 15, 37] developed by different groups. Eqs. (1) and (2) are intended to be the general formulas based on the basic ideas of the combination mechanism. The different models are examples of the general case that we considered in these equations. In these models, different method(s) and/or assumption(s) are usually introduced to construct the precise form of the kernel functions  $\mathcal{R}_{M_j, q_1 \bar{q}_2}(p_{M_j}, p_1, p_2; N_{q_i}, N_{\bar{q}_i})$  and  $\mathcal{R}_{B_j, q_1 q_2 q_3}(p_{B_j}, p_1, p_2, p_3; N_{q_i}, N_{\bar{q}_i})$  in order to provide a good description of different properties of the hadrons, such as momentum distributions and so on. For example, in the recombination model developed by Hwa and collaborators [12], these kernel functions are just the recombination functions.

The joint distributions  $f_{q_1 \bar{q}_2}$  and  $f_{q_1 q_2 q_3}$  are the number densities that satisfy

$$\int f_{q_1 \bar{q}_2}(p_1, p_2; N_{q_i}, N_{\bar{q}_i}) dp_1 dp_2 = N_{q_1 \bar{q}_2}, \quad (3)$$

$$\int f_{q_1 q_2 q_3}(p_1, p_2, p_3; N_{q_i}, N_{\bar{q}_i}) dp_1 dp_2 dp_3 = N_{q_1 q_2 q_3}, \quad (4)$$

respectively, where  $N_{q_1 \bar{q}_2} = N_{q_1} N_{\bar{q}_2}$ , and

$$N_{q_1 q_2 q_3} = \begin{cases} N_{q_1} N_{q_2} N_{q_3} & \text{for } q_1 \neq q_2 \neq q_3 \\ N_{q_1} (N_{q_1} - 1) N_{q_3} & \text{for } q_1 = q_2 \neq q_3 \\ N_{q_1} (N_{q_1} - 1) (N_{q_1} - 2) & \text{for } q_1 = q_2 = q_3 \end{cases} \quad (5)$$

are the numbers of all the possible  $(q\bar{q})$ 's and  $(qqq)$ 's in the bulk quark-antiquark system that we consider. For the convenience of comparison, we rewrite them as

$$f_{q_1 \bar{q}_2}(p_1, p_2; N_{q_i}, N_{\bar{q}_i}) = N_{q_1 \bar{q}_2} f_{q_1 \bar{q}_2}^{(n)}(p_1, p_2; N_{q_i}, N_{\bar{q}_i}), \quad (6)$$

$$f_{q_1 q_2 q_3}(p_1, p_2, p_3; N_{q_i}, N_{\bar{q}_i}) = N_{q_1 q_2 q_3} f_{q_1 q_2 q_3}^{(n)}(p_1, p_2, p_3; N_{q_i}, N_{\bar{q}_i}), \quad (7)$$

so that the distributions are normalized to unity where we denote by using the superscript  $(n)$ , i.e.,

$$\int f_{q_1 \bar{q}_2}^{(n)}(p_1, p_2; N_{q_i}, N_{\bar{q}_i}) dp_1 dp_2 = 1, \quad (8)$$

$$\int f_{q_1 q_2 q_3}^{(n)}(p_1, p_2, p_3; N_{q_i}, N_{\bar{q}_i}) dp_1 dp_2 dp_3 = 1. \quad (9)$$

In terms of these normalized joint distributions, we have

$$f_{M_j}(p_{M_j}) = \sum_{q_1 \bar{q}_2} N_{q_1 \bar{q}_2} \int dp_1 dp_2 f_{q_1 \bar{q}_2}^{(n)}(p_1, p_2; N_{q_i}, N_{\bar{q}_i}) \times \mathcal{R}_{M_j}(p_{M_j}, p_1, p_2; N_{q_i}, N_{\bar{q}_i}), \quad (10)$$

$$f_{B_j}(p_{B_j}) = \sum_{q_1 q_2 q_3} N_{q_1 q_2 q_3} \int dp_1 dp_2 dp_3 \times f_{q_1 q_2 q_3}^{(n)}(p_1, p_2, p_3; N_{q_i}, N_{\bar{q}_i}) \times \mathcal{R}_{B_j}(p_{B_j}, p_1, p_2, p_3; N_{q_i}, N_{\bar{q}_i}). \quad (11)$$

Integrating over  $p_{M_j}$  or  $p_{B_j}$  from the momentum distributions, we obtain the average numbers of the directly produced mesons  $M_j$  and baryons  $B_j$  as

$$\bar{N}_{M_j}(N_{q_i}, N_{\bar{q}_i}) = \sum_{q_1 \bar{q}_2} N_{q_1 \bar{q}_2} \int dp_{M_j} dp_1 dp_2 \times f_{q_1 \bar{q}_2}^{(n)}(p_1, p_2; N_{q_i}, N_{\bar{q}_i}) \mathcal{R}_{M_j}(p_{M_j}, p_1, p_2; N_{q_i}, N_{\bar{q}_i}), \quad (12)$$

$$\bar{N}_{B_j}(N_{q_i}, N_{\bar{q}_i}) = \sum_{q_1 q_2 q_3} N_{q_1 q_2 q_3} \int dp_{B_j} dp_1 dp_2 dp_3 \times f_{q_1 q_2 q_3}^{(n)}(p_1, p_2, p_3; N_{q_i}, N_{\bar{q}_i}) \mathcal{R}_{B_j}(p_{B_j}, p_1, p_2, p_3; N_{q_i}, N_{\bar{q}_i}). \quad (13)$$

For a reaction at a given energy, the average numbers of quarks,  $\langle N_{q_i} \rangle$ , and those for the antiquarks,  $\langle N_{\bar{q}_i} \rangle$ , of different flavors are fixed. The numbers of quarks and antiquarks follow a certain distribution which we denote by  $P(N_{q_i}, N_{\bar{q}_i}, \langle N_{q_i} \rangle, \langle N_{\bar{q}_i} \rangle)$ . The average yields of mesons and baryons are given by

$$\langle N_{M_j} \rangle(\sqrt{s}) = \sum_{N_{q_i} N_{\bar{q}_i}} P(N_{q_i}, N_{\bar{q}_i}, \langle N_{q_i} \rangle, \langle N_{\bar{q}_i} \rangle) \bar{N}_{M_j}(N_{q_i}, N_{\bar{q}_i}),$$

$$\langle N_{B_j} \rangle(\sqrt{s}) = \sum_{N_{q_i} N_{\bar{q}_i}} P(N_{q_i}, N_{\bar{q}_i}, \langle N_{q_i} \rangle, \langle N_{\bar{q}_i} \rangle) \bar{N}_{B_j}(N_{q_i}, N_{\bar{q}_i}).$$

These equations are the general formalism for calculating the average yield of a certain sort of hadrons in high-energy reactions based on the basic ideas of the combination mechanism. More specific results can be obtained for special cases when special assumptions are made about the distributions and/or the kernel functions. We present such cases step by step in the following.

## B. Factorization of flavor and momentum dependences

The flavor dependence of the kernel functions  $\mathcal{R}_{M_j, q_1 \bar{q}_2}$  and  $\mathcal{R}_{B_j, q_1 q_2 q_3}$  is responsible for flavor conservation in the combination process and the differences between the combination probabilities for different flavors of quarks, antiquarks, and hadrons. In general, the momentum and the flavor dependencies of these kernel functions are coupled to each other. In that case, the results for the ratios of the average yields of different hadrons can be dependent on the way of coupling. In this paper, we do not consider such coupling effects. In contrast, in the following, we consider only the simplest case where the momentum and flavor dependencies of the kernel functions are decoupled from each other. In other words, we consider the case where they are factorized, i.e.,

$$\mathcal{R}_{M_j, q_1 \bar{q}_2}(p_M, p_1, p_2; N_q, N_{\bar{q}}) = \mathcal{R}_{M_j, q_1 \bar{q}_2}^{(f)}(N_q, N_{\bar{q}}) \mathcal{R}_M^{(p)}(p_M, p_1, p_2; N_q, N_{\bar{q}}), \quad (14)$$

$$\mathcal{R}_{B_j, q_1 q_2 q_3}(p_B, p_1, p_2, p_3; N_q, N_{\bar{q}}) = \mathcal{R}_{B_j, q_1 q_2 q_3}^{(f)}(N_q, N_{\bar{q}}) \mathcal{R}_B^{(p)}(p_B, p_1, p_2, p_3; N_q, N_{\bar{q}}), \quad (15)$$

where the flavor-independent parts  $\mathcal{R}_M^{(p)}(p_M, p_1, p_2; N_q, N_{\bar{q}})$  and  $\mathcal{R}_B^{(p)}(p_B, p_1, p_2, p_3; N_q, N_{\bar{q}})$  denote the probability for a  $(q\bar{q})$  with momenta  $p_1$  and  $p_2$  in a system consisting of  $N_q$  quarks and  $N_{\bar{q}}$  antiquarks to combine with each other to form a meson  $M$  with momentum  $p_M$  and that for a  $(qqq)$  with momenta  $p_1, p_2$ , and  $p_3$  in the system to combine with each other to form a baryon  $B$  with momentum  $p_B$ , respectively. The flavor-dependent parts  $\mathcal{R}_{M_j, q_1 \bar{q}_2}^{(f)}(N_q, N_{\bar{q}})$  and  $\mathcal{R}_{B_j, q_1 q_2 q_3}^{(f)}(N_q, N_{\bar{q}})$  represent the probability for the  $q_1$  and  $\bar{q}_2$  to combine into the specified meson  $M_j$  in the case that they are known to combine into a meson and that for the  $q_1, q_2$ , and  $q_3$  to combine into the specified baryon  $B_j$  in the case that they are known to combine into a baryon, respectively. They are taken as satisfying the normalization condition

$$\sum_j \mathcal{R}_{M_j, q_1 \bar{q}_2}^{(f)} = 1, \quad (16)$$

$$\sum_j \mathcal{R}_{B_j, q_1 q_2 q_3}^{(f)} = 1. \quad (17)$$

We further assume that the normalized joint momentum distributions of the quarks and/or antiquarks are flavor in-

dependent, i.e.,

$$f_{q_1 \bar{q}_2}^{(n)}(p_1, p_2; N_q, N_{\bar{q}}) = f_{q\bar{q}}^{(n)}(p_1, p_2; N_q, N_{\bar{q}}), \quad (18)$$

$$f_{q_1 q_2 q_3}^{(n)}(p_1, p_2, p_3; N_q, N_{\bar{q}}) = f_{qqq}^{(n)}(p_1, p_2, p_3; N_q, N_{\bar{q}}). \quad (19)$$

Under these two approximations, we have

$$\begin{aligned} \bar{N}_{M_j}(N_q, N_{\bar{q}}) &= \sum_{q_1 \bar{q}_2} N_{q_1 \bar{q}_2} \mathcal{R}_{M_j, q_1 \bar{q}_2}^{(f)}(N_q, N_{\bar{q}}) \times \\ &\quad \int dp_M dp_1 dp_2 f_{q\bar{q}}^{(n)}(p_1, p_2; N_q, N_{\bar{q}}) \times \\ &\quad \mathcal{R}_M^{(p)}(p_M, p_1, p_2; N_q, N_{\bar{q}}), \end{aligned} \quad (20)$$

$$\begin{aligned} \bar{N}_{B_j}(N_q, N_{\bar{q}}) &= \sum_{q_1 q_2 q_3} N_{q_1 q_2 q_3} \mathcal{R}_{B_j, q_1 q_2 q_3}^{(f)}(N_q, N_{\bar{q}}) \times \\ &\quad \int dp_B dp_1 dp_2 dp_3 f_{qqq}^{(n)}(p_1, p_2, p_3; N_q, N_{\bar{q}}) \times \\ &\quad \mathcal{R}_B^{(p)}(p_B, p_1, p_2, p_3; N_q, N_{\bar{q}}). \end{aligned} \quad (21)$$

We denote

$$\gamma_M(N_q, N_{\bar{q}}, \sqrt{s}) = \int dp_M dp_1 dp_2 \times f_{q\bar{q}}^{(n)}(p_1, p_2; N_q, N_{\bar{q}}) \mathcal{R}_M^{(p)}(p_M, p_1, p_2; N_q, N_{\bar{q}}), \quad (22)$$

$$\gamma_B(N_q, N_{\bar{q}}, \sqrt{s}) = \int dp_B dp_1 dp_2 dp_3 \times f_{qqq}^{(n)}(p_1, p_2, p_3; N_q, N_{\bar{q}}) \mathcal{R}_B^{(p)}(p_B, p_1, p_2, p_3; N_q, N_{\bar{q}}), \quad (23)$$

and obtain

$$\bar{N}_{M_j}(N_q, N_{\bar{q}}) = \sum_{q_1 \bar{q}_2} N_{q_1 \bar{q}_2} \mathcal{R}_{M_j, q_1 \bar{q}_2}^{(f)} \gamma_M, \quad (24)$$

$$\bar{N}_{B_j}(N_q, N_{\bar{q}}) = \sum_{q_1 q_2 q_3} N_{q_1 q_2 q_3} \mathcal{R}_{B_j, q_1 q_2 q_3}^{(f)} \gamma_B. \quad (25)$$

Summing over different species of mesons and those of the baryons, respectively, we obtain the average total numbers of mesons and baryons produced in the combination of the system of  $N_q$  quarks and  $N_{\bar{q}}$  antiquarks as

$$\bar{N}_M(N_q, N_{\bar{q}}, \sqrt{s}) = N_{q\bar{q}} \gamma_M(N_q, N_{\bar{q}}, \sqrt{s}), \quad (26)$$

$$\bar{N}_B(N_q, N_{\bar{q}}, \sqrt{s}) = N_{qqq} \gamma_B(N_q, N_{\bar{q}}, \sqrt{s}), \quad (27)$$

where  $N_{q\bar{q}} = N_q N_{\bar{q}}$  and  $N_{qqq} = N_q(N_q - 1)(N_q - 2)$  are the total number of  $q\bar{q}$  pairs and that of  $qqq$  systems, respectively. The factors  $\gamma_M(N_q, N_{\bar{q}}, \sqrt{s})$  and  $\gamma_B(N_q, N_{\bar{q}}, \sqrt{s})$  represent the probability for a particular  $q\bar{q}$  from the system consisting of  $N_q$  quarks and  $N_{\bar{q}}$  antiquarks to combine with each other to form a meson and that for a  $qqq$  to form a baryon, respectively. We emphasize in particular that Eq. (26) does not mean that the average yield of mesons is proportional to the product of the number of quarks and that of antiquarks since the factor  $\gamma_M$  can depend strongly on  $N_q$  and  $N_{\bar{q}}$ . This is because, for a given  $q$ , the larger  $N_q$  and/or  $N_{\bar{q}}$ , the more possibilities for the  $q$  to combine with the others to form a hadron, and thus the smaller the

probability for it to combine with the given  $\bar{q}$  to form the meson. In fact, it can in general be expected that  $\gamma_M$  should be more or less inversely proportional to  $N_q$  and/or  $N_{\bar{q}}$  and the final result for  $\langle N_M \rangle$  should be roughly proportional to  $N_q + N_{\bar{q}}$ . A similar conclusion holds for  $\langle N_B \rangle$ .

In terms of the total average numbers of mesons and baryons, the average number of a specified meson  $M_j$  and that of a specified baryon  $B_j$  are given by

$$\bar{N}_{M_j}(N_{q_i}, N_{\bar{q}_i}) = \sum_{q_1 \bar{q}_2} \frac{N_{q_1 \bar{q}_2}}{N_{q\bar{q}}} \mathcal{R}_{M_j, q_1 \bar{q}_2}^{(f)}(N_{q_i}, N_{\bar{q}_i}) \times \bar{N}_M(N_q, N_{\bar{q}}, \sqrt{s}), \quad (28)$$

$$\bar{N}_{B_j}(N_{q_i}, N_{\bar{q}_i}) = \sum_{q_1 q_2 q_3} \frac{N_{q_1 q_2 q_3}}{N_{qqq}} \mathcal{R}_{B_j, q_1 q_2 q_3}^{(f)}(N_{q_i}, N_{\bar{q}_i}) \times \bar{N}_B(N_q, N_{\bar{q}}, \sqrt{s}). \quad (29)$$

The parts describing the flavor dependence of the kernel functions,  $\mathcal{R}_{M_j, q_1 \bar{q}_2}^{(f)}(N_{q_i}, N_{\bar{q}_i})$  and  $\mathcal{R}_{B_j, q_1 q_2 q_3}^{(f)}(N_{q_i}, N_{\bar{q}_i})$ , have to guarantee flavor conservation in the combination process. Hence, they contain the Kronecker  $\delta$ 's and constant factors  $C_{M_j}$  and  $C_{B_j}$ . For example, for  $\pi^+$  and  $p$ , they are given by

$$\begin{aligned} \mathcal{R}_{\pi^+, q_1 \bar{q}_2}^{(f)} &= C_{\pi^+} \delta_{q_1, u} \delta_{\bar{q}_2, \bar{d}}, \\ \mathcal{R}_{p, q_1 q_2 q_3}^{(f)} &= C_p (\delta_{q_1, u} \delta_{q_2, u} \delta_{q_3, d} + \delta_{q_1, u} \delta_{q_2, d} \delta_{q_3, u} \\ &\quad + \delta_{q_1, d} \delta_{q_2, u} \delta_{q_3, u}). \end{aligned} \quad (30)$$

We recall that, in the factorized case considered here, the flavor-dependent part  $\mathcal{R}_{M_j, q_1 \bar{q}_2}^{(f)}(N_{q_i}, N_{\bar{q}_i})$  of the kernel function represents the probability for the specified  $q_1 \bar{q}_2$  with the specified flavor  $q_1$  and  $\bar{q}_2$  from the system consisting of  $N_q$  quarks and  $N_{\bar{q}}$  antiquarks to form the specified meson  $M_j$  under the condition that they are known to form a meson. Although we can not prove it, it is very unlikely that this probability still depends strongly on the environment. We therefore consider the simplified case where  $C_{M_j}$  is taken as a constant independent of  $N_q$  or  $N_{\bar{q}}$ . The same applies to  $\mathcal{R}_{B_j, q_1 q_2 q_3}^{(f)}(N_{q_i}, N_{\bar{q}_i})$ . In this case, we have

$$\bar{N}_{M_j}(N_{q_i}, N_{\bar{q}_i}) = C_{M_j} \frac{N_{q_1 \bar{q}_2}}{N_{q\bar{q}}} \bar{N}_M(N_q, N_{\bar{q}}, \sqrt{s}), \quad (31)$$

$$\bar{N}_{B_j}(N_{q_i}, N_{\bar{q}_i}) = N_{iter} C_{B_j} \frac{N_{q_1 q_2 q_3}}{N_{qqq}} \bar{N}_B(N_q, N_{\bar{q}}, \sqrt{s}), \quad (32)$$

where  $N_{iter}$  stands for the number of possible iterations of  $q_1 q_2 q_3$  which is 1, 3, and 6 for three identical flavor, two different flavor, and three different flavor cases, respectively.

In the case when only  $J^P = 0^-$  and  $1^-$  mesons and  $J^P = \frac{1}{2}^+$  and  $\frac{3}{2}^+$  baryons are considered, we have, for mesons,

$$C_{M_j} = \begin{cases} 1/(1 + R_{V/P}) & \text{for } J^P = 0^- \text{ mesons,} \\ R_{V/P}/(1 + R_{V/P}) & \text{for } J^P = 1^- \text{ mesons,} \end{cases} \quad (33)$$

where  $R_{V/P}$  represents the ratio of the  $J^P = 1^-$  vector mesons to the  $J^P = 0^-$  pseudoscalar mesons of the same

flavor composition; and for baryons,

$$C_{B_j} = \begin{cases} R_{O/D}/(1 + R_{O/D}) & \text{for } J^P = (1/2)^+ \text{ baryons,} \\ 1/(1 + R_{O/D}) & \text{for } J^P = (3/2)^+ \text{ baryons,} \end{cases} \quad (34)$$

except that  $C_{\Lambda} = C_{\Sigma^0} = R_{O/D}/(1 + 2R_{O/D})$ ,  $C_{\Sigma^{*0}} = 1/(1 + 2R_{O/D})$ , and  $C_{\Delta^{++}} = C_{\Delta^-} = C_{\Omega^-} = 1$ . Here,  $R_{O/D}$  stands for the ratio of the  $J^P = (1/2)^+$  octet to the  $J^P = (3/2)^+$  decuplet baryons of the same flavor composition. The two parameters  $R_{V/P}$  and  $R_{O/D}$  can be determined by using the data from different high-energy reactions [1, 13, 38].

### C. Modeling $P(N_{q_i}, \langle N_{q_i} \rangle, \sqrt{s})$

We consider three flavors  $u$ ,  $d$ , and  $s$  of quarks and antiquarks. Inside the system of  $N_q$  quarks and  $N_{\bar{q}}$  antiquarks, we suppose that each quark can take flavor  $u$ ,  $d$ , or  $s$  with given probability  $p_u$ ,  $p_d$ , or  $p_s$  independent of the others. In this case, the numbers of  $u$ ,  $d$ , and  $s$  quarks inside the system at a given  $N_q$  obey the multinomial distribution, i.e.,

$$B(N_{q_i}; N_q) = \frac{N_q!}{N_u! N_d! N_s!} p_u^{N_u} p_d^{N_d} p_s^{N_s} \delta_{N_q, N_u + N_d + N_s}, \quad (35)$$

where  $p_u = p_d = 1/(2 + \lambda_q)$ ,  $p_s = \lambda_q/(2 + \lambda_q)$ , and  $\lambda_q$  is the effective strangeness suppression factor for quarks. Similarly, for the antiquarks, we have

$$B(N_{\bar{q}_i}; N_{\bar{q}}) = \frac{N_{\bar{q}}!}{N_{\bar{u}}! N_{\bar{d}}! N_{\bar{s}}!} p_{\bar{u}}^{N_{\bar{u}}} p_{\bar{d}}^{N_{\bar{d}}} p_{\bar{s}}^{N_{\bar{s}}} \delta_{N_{\bar{q}}, N_{\bar{u}} + N_{\bar{d}} + N_{\bar{s}}}, \quad (36)$$

where  $p_{\bar{u}} = p_{\bar{d}} = 1/(2 + \lambda)$ ,  $p_{\bar{s}} = \lambda/(2 + \lambda)$ , and  $\lambda$  is the strangeness suppression factor for antiquark production.

In general, in high-energy heavy-ion collisions, the system contains the contributions of the net quarks coming from the incident nuclei. Hence the effective strangeness suppression factor  $\lambda_q$  for the quarks is different from  $\lambda$  for the antiquarks, which do not have influence from the net quarks. Here, we keep them as distinguished from each other so that we can apply the results to different cases. Furthermore, we emphasize that the system considered corresponds to a quark-antiquark system produced in an AA collision in a limited kinematic region. The system is supposed to be a small part of the whole quark-antiquark system produced in the reaction so that the influence from the global flavor compensation is considered to be negligible. The global flavor compensation can have some influence on the flavor correlation in hadron production. Such a case was for example discussed in [5] for baryon-antibaryon flavor correlation in  $e^+e^-$  annihilation where the number of quarks was of the order of tens and the average yield of baryons in an event was less than one. It was found that, even in that case, the global flavor compensation does have some effect on the flavor correlation but the effect is not very large. Hence, for simplicity and clarity, we neglect them in the discussion here.



At given  $N_q$  and  $N_{\bar{q}}$ , we average over the distribution of the numbers of quarks and/or antiquarks for different flavors. It can easily be shown that, for  $q_1 \neq q_2 \neq q_3$ ,

$$\sum_{N_{q_i}} N_{q_i} B(N_{q_i}; N_q) = N_q p_{q_1}, \quad (37)$$

$$\sum_{N_{q_i}} N_{q_1} N_{q_2} N_{q_3} B(N_{q_i}; N_q) = N_{qqq} p_{q_1} p_{q_2} p_{q_3}, \quad (38)$$

$$\sum_{N_{q_i}} N_{q_1} (N_{q_1} - 1) N_{q_2} B(N_{q_i}; N_q) = N_{qqq} p_{q_1}^2 p_{q_2}, \quad (39)$$

and similarly for others, so we obtain

$$\bar{N}_{M_j}(N_q, N_{\bar{q}}, \sqrt{s}) = C_{M_j} p_{q_1} p_{\bar{q}_2} \bar{N}_M(N_q, N_{\bar{q}}, \sqrt{s}), \quad (40)$$

$$\bar{N}_{B_j}(N_q, N_{\bar{q}}, \sqrt{s}) = N_{iter} C_{B_j} p_{q_1} p_{q_2} p_{q_3} \bar{N}_B(N_q, N_{\bar{q}}, \sqrt{s}). \quad (41)$$

For a subsystem of quarks and antiquarks in a given kinematic region in AA collisions at given energy  $\sqrt{s}$ ,  $\langle N_q \rangle$  and  $\langle N_{\bar{q}} \rangle$  are fixed while  $N_q$  and  $N_{\bar{q}}$  follow the distributions  $P_q(N_q; \langle N_q \rangle)$  and  $P_{\bar{q}}(N_{\bar{q}}; \langle N_{\bar{q}} \rangle)$ , respectively. Hence, we need to average over these distributions and obtain

$$\langle N_{M_j} \rangle(\langle N_q \rangle, \langle N_{\bar{q}} \rangle, \sqrt{s}) = C_{M_j} p_{q_1} p_{\bar{q}_2} \langle N_M \rangle, \quad (42)$$

$$\langle N_{B_j} \rangle(\langle N_q \rangle, \langle N_{\bar{q}} \rangle, \sqrt{s}) = N_{iter} C_{B_j} p_{q_1} p_{q_2} p_{q_3} \langle N_B \rangle, \quad (43)$$

where  $\langle N_M \rangle$  and  $\langle N_B \rangle$  are functions of  $\langle N_q \rangle$ ,  $\langle N_{\bar{q}} \rangle$  and  $\sqrt{s}$  and stand for the average total number of the mesons and that of the baryons produced in the combination process. They are given by

$$\langle N_M \rangle = \sum_{N_q N_{\bar{q}}} P_q(N_q; \langle N_q \rangle) P_{\bar{q}}(N_{\bar{q}}; \langle N_{\bar{q}} \rangle) \bar{N}_M(N_q, N_{\bar{q}}, \sqrt{s}),$$

$$\langle N_B \rangle = \sum_{N_q N_{\bar{q}}} P_q(N_q; \langle N_q \rangle) P_{\bar{q}}(N_{\bar{q}}; \langle N_{\bar{q}} \rangle) \bar{N}_B(N_q, N_{\bar{q}}, \sqrt{s}).$$

We see that, in this case, for the directly produced hadrons, the ratios of the yields of different mesons, those of different baryons and those of the antibaryons separately are constants depending on the parameters  $\lambda$ ,  $\lambda_q$ ,  $R_{V/P}$  and  $R_{O/D}$ . In general the effective strangeness suppression factor  $\lambda_q$  for quarks contains the influence from the net quark contributions and can be dependent on  $\langle N_q \rangle$  and  $\langle N_{\bar{q}} \rangle$ . This leads to a dependence on  $\langle N_q \rangle$  and  $\langle N_{\bar{q}} \rangle$  even for such kinds of hadron yield ratios. In the case that the net quark contribution is negligible, we have  $\lambda_q = \lambda$ ; these kinds of hadron yield ratios become constants independent of  $\langle N_q \rangle$  and  $\langle N_{\bar{q}} \rangle$ . This should be the case for a subsample in the central rapidity region of the bulk quark-antiquark system produced in AA collisions at very high energies such as those at the LHC. In this case, we have also  $\langle N_B \rangle = \langle N_{\bar{B}} \rangle$  and this, together with  $p_{q_i} = p_{\bar{q}_i}$ , leads to  $\langle N_{B_j} \rangle = \langle N_{\bar{B}_j} \rangle$ . These are predictions that can be checked at the LHC.

#### D. Including the decay contributions

Including the decay contributions, we calculate the yields of different hadrons in the final state. We denote the

decay contribution from a hadron  $h_i$  to  $h_j$  by  $Br(h_i \rightarrow h_j)$  and obtain

$$\langle N_{h_j}^f \rangle = \langle N_{h_j} \rangle + \sum_{i \neq j} Br(h_i \rightarrow h_j) \langle N_{h_i} \rangle, \quad (44)$$

where we use the superscript  $f$  to denote the results for the final hadrons to differentiate them from those for the directly produced hadrons. Here we consider only the influence from the decay of the short-lived hadrons but do not consider the influences from the final-state interactions of the hadrons.

The value of  $Br(h_i \rightarrow h_j)$  can be obtained easily from the materials given by the Particle Data Group [39]. In the following, we take the strong and the electromagnetic decays into account. For most of the hadrons, the results look very simple. In the case in which only  $J^P = 0^-$  and  $1^-$  mesons and  $J^P = (1/2)^+$  and  $(3/2)^+$  baryons are included, the average yields of final hadrons, e.g.,  $K^+$ ,  $p$ , and  $\Lambda$ , are given as

$$\langle N_{K^+}^f \rangle = \langle N_{K^+} \rangle + \frac{2}{3} \langle N_{K^{*0}} \rangle + \frac{1}{3} \langle N_{K^{*+}} \rangle + 0.489 \langle N_{\phi} \rangle, \quad (45)$$

$$\langle N_p^f \rangle = \langle N_p \rangle + \langle N_{\Delta^{++}} \rangle + \frac{2}{3} \langle N_{\Delta^+} \rangle + \frac{1}{3} \langle N_{\Delta^0} \rangle, \quad (46)$$

$$\langle N_{\Lambda}^f \rangle = \langle N_{\Lambda} \rangle + \langle N_{\Sigma^0} \rangle + 0.883 \langle N_{\Sigma^{*0}} \rangle + 0.94 (\langle N_{\Sigma^{*+}} \rangle + \langle N_{\Sigma^{*-}} \rangle). \quad (47)$$

We consider the case discussed in Sec. IIC and substitute the results for  $\langle N_{M_j} \rangle$  and those for  $\langle N_{B_j} \rangle$  given by Eqs. (42) and (43) into the above equations and obtain

$$\langle N_{K^+}^f \rangle = p_u p_{\bar{s}} \langle N_M \rangle + \frac{0.489 R_{V/P}}{1 + R_{V/P}} p_s p_{\bar{s}} \langle N_M \rangle, \quad (48)$$

$$\langle N_p^f \rangle = 4 p_u^3 \langle N_B \rangle, \quad (49)$$

$$\langle N_{\Lambda}^f \rangle = \left( \frac{5.30 + 12 R_{O/D}}{2 R_{O/D} + 1} + \frac{5.64}{R_{O/D} + 1} \right) p_u^2 p_s \langle N_B \rangle, \quad (50)$$

where we have taken  $p_u = p_d$  and  $p_{\bar{u}} = p_{\bar{d}}$ .

Taking  $R_{V/P} = 3$  according to the spin counting, but  $R_{O/D} = 2$  since decuplet baryon production is observed much more suppressed [1, 13, 38], we then have

$$\langle N_{K^+}^f \rangle = (p_u p_{\bar{s}} + 0.37 p_s p_{\bar{s}}) \langle N_M \rangle, \quad (51)$$

$$\langle N_{\Lambda}^f \rangle = 7.74 p_u^2 p_s \langle N_B \rangle. \quad (52)$$

In Table I, we show the results obtained for baryons and antibaryons and those for strange mesons in different cases. We see that, like for protons, the results for many final baryons look even simpler than those for the directly produced cases since the corresponding decuplet baryons decay strongly to these baryons. This makes the average final yields for these baryons independent of the ratio  $R_{O/D}$ . In Table I, we also present the results for the simple case without net quarks. The results obtained if we take  $R_{V/P} = 3$  and  $R_{O/D} = 2$  are also given.

TABLE I: Average yields of the directly produced hadrons and those including decay contributions. Here, in the second column, we show the results for the directly produced hadrons. In the third column, we show the results when the strong and the electromagnetic (S & EM) decay contributions are taken into account. The fourth column shows the results when the net quark influence is negligible. In the last two columns, we see the results for the case when  $R_{V/P} = 3$  and  $R_{O/D} = 2$ .

Hadron	Directly produced	With S & EM decays	$N_q^{net} = 0$ ( $\lambda_q = \lambda$ )	$R_{V/P} = 3, R_{O/D} = 2$	
				$N_q^{net} \neq 0$	$N_q^{net} = 0$
$p$	$\frac{3R_{O/D}}{1+R_{O/D}} p_u^3 \langle N_B \rangle$	$4p_u^3 \langle N_B \rangle$	$\frac{4}{(2+\lambda)^3} \langle N_B \rangle$	$4p_u^3 \langle N_B \rangle$	$\frac{4}{(2+\lambda)^3} \langle N_B \rangle$
$n$	$\frac{3R_{O/D}}{1+R_{O/D}} p_u^3 \langle N_B \rangle$	$4p_u^3 \langle N_B \rangle$	$\frac{4}{(2+\lambda)^3} \langle N_B \rangle$	$4p_u^3 \langle N_B \rangle$	$\frac{4}{(2+\lambda)^3} \langle N_B \rangle$
$\Xi^0$	$\frac{3R_{O/D}}{1+R_{O/D}} p_u p_s^2 \langle N_B \rangle$	$3p_u p_s^2 \langle N_B \rangle$	$\frac{3\lambda^2}{(2+\lambda)^3} \langle N_B \rangle$	$3p_u p_s^2 \langle N_B \rangle$	$\frac{3\lambda^2}{(2+\lambda)^3} \langle N_B \rangle$
$\Xi^-$	$\frac{3R_{O/D}}{1+R_{O/D}} p_u p_s^2 \langle N_B \rangle$	$3p_u p_s^2 \langle N_B \rangle$	$\frac{3\lambda^2}{(2+\lambda)^3} \langle N_B \rangle$	$3p_u p_s^2 \langle N_B \rangle$	$\frac{3\lambda^2}{(2+\lambda)^3} \langle N_B \rangle$
$\Omega^-$	$p_s^3 \langle N_B \rangle$	$p_s^3 \langle N_B \rangle$	$\frac{\lambda^3}{(2+\lambda)^3} \langle N_B \rangle$	$p_s^3 \langle N_B \rangle$	$\frac{\lambda^3}{(2+\lambda)^3} \langle N_B \rangle$
$\bar{p}$	$\frac{3R_{O/D}}{1+R_{O/D}} p_{\bar{u}}^3 \langle N_{\bar{B}} \rangle$	$4p_{\bar{u}}^3 \langle N_{\bar{B}} \rangle$	$\frac{4}{(2+\lambda)^3} \langle N_B \rangle$	$4p_{\bar{u}}^3 \langle N_{\bar{B}} \rangle$	$\frac{4}{(2+\lambda)^3} \langle N_B \rangle$
$\bar{n}$	$\frac{3R_{O/D}}{1+R_{O/D}} p_{\bar{u}}^3 \langle N_{\bar{B}} \rangle$	$4p_{\bar{u}}^3 \langle N_{\bar{B}} \rangle$	$\frac{4}{(2+\lambda)^3} \langle N_B \rangle$	$4p_{\bar{u}}^3 \langle N_{\bar{B}} \rangle$	$\frac{4}{(2+\lambda)^3} \langle N_B \rangle$
$\bar{\Xi}^0$	$\frac{3R_{O/D}}{1+R_{O/D}} p_{\bar{u}} p_{\bar{s}}^2 \langle N_{\bar{B}} \rangle$	$3p_{\bar{u}} p_{\bar{s}}^2 \langle N_{\bar{B}} \rangle$	$\frac{3\lambda^2}{(2+\lambda)^3} \langle N_B \rangle$	$3p_{\bar{u}} p_{\bar{s}}^2 \langle N_{\bar{B}} \rangle$	$\frac{3\lambda^2}{(2+\lambda)^3} \langle N_B \rangle$
$\bar{\Xi}^+$	$\frac{3R_{O/D}}{1+R_{O/D}} p_{\bar{u}} p_{\bar{s}}^2 \langle N_{\bar{B}} \rangle$	$3p_{\bar{u}} p_{\bar{s}}^2 \langle N_{\bar{B}} \rangle$	$\frac{3\lambda^2}{(2+\lambda)^3} \langle N_B \rangle$	$3p_{\bar{u}} p_{\bar{s}}^2 \langle N_{\bar{B}} \rangle$	$\frac{3\lambda^2}{(2+\lambda)^3} \langle N_B \rangle$
$\bar{\Omega}^+$	$p_{\bar{s}}^3 \langle N_{\bar{B}} \rangle$	$p_{\bar{s}}^3 \langle N_{\bar{B}} \rangle$	$\frac{\lambda^3}{(2+\lambda)^3} \langle N_B \rangle$	$p_{\bar{s}}^3 \langle N_{\bar{B}} \rangle$	$\frac{\lambda^3}{(2+\lambda)^3} \langle N_B \rangle$
$K^+$	$\frac{1}{1+R_{V/P}} p_u p_{\bar{s}} \langle N_M \rangle$	$p_u p_{\bar{s}} (1 + \frac{0.49R_{V/P}}{1+R_{V/P}} \lambda_q) \langle N_M \rangle$	$\frac{\lambda}{(2+\lambda)^2} (1 + \frac{0.49R_{V/P}}{1+R_{V/P}} \lambda) \langle N_M \rangle$	$p_u p_{\bar{s}} (1 + 0.37\lambda_q) \langle N_M \rangle$	$\frac{\lambda+0.37\lambda^2}{(2+\lambda)^2} \langle N_M \rangle$
$K^-$	$\frac{1}{1+R_{V/P}} p_{\bar{u}} p_s \langle N_M \rangle$	$p_{\bar{u}} p_s (1 + \frac{0.49R_{V/P}}{1+R_{V/P}} \lambda) \langle N_M \rangle$	$\frac{\lambda}{(2+\lambda)^2} (1 + \frac{0.49R_{V/P}}{1+R_{V/P}} \lambda) \langle N_M \rangle$	$p_{\bar{u}} p_s (1 + 0.37\lambda) \langle N_M \rangle$	$\frac{\lambda+0.37\lambda^2}{(2+\lambda)^2} \langle N_M \rangle$
$K^0$	$\frac{1}{1+R_{V/P}} p_u p_{\bar{s}} \langle N_M \rangle$	$p_u p_{\bar{s}} (1 + \frac{0.34R_{V/P}}{1+R_{V/P}} \lambda_q) \langle N_M \rangle$	$\frac{\lambda}{(2+\lambda)^2} (1 + \frac{0.34R_{V/P}}{1+R_{V/P}} \lambda) \langle N_M \rangle$	$p_u p_{\bar{s}} (1 + 0.26\lambda_q) \langle N_M \rangle$	$\frac{\lambda+0.26\lambda^2}{(2+\lambda)^2} \langle N_M \rangle$
$\bar{K}^0$	$\frac{1}{1+R_{V/P}} p_{\bar{u}} p_s \langle N_M \rangle$	$p_{\bar{u}} p_s (1 + \frac{0.34R_{V/P}}{1+R_{V/P}} \lambda) \langle N_M \rangle$	$\frac{\lambda}{(2+\lambda)^2} (1 + \frac{0.34R_{V/P}}{1+R_{V/P}} \lambda) \langle N_M \rangle$	$p_{\bar{u}} p_s (1 + 0.26\lambda) \langle N_M \rangle$	$\frac{\lambda+0.26\lambda^2}{(2+\lambda)^2} \langle N_M \rangle$
$\phi$	$\frac{R_{V/P}}{1+R_{V/P}} p_s p_{\bar{s}} \langle N_M \rangle$	$\frac{R_{V/P}}{1+R_{V/P}} p_s p_{\bar{s}} \langle N_M \rangle$	$(\frac{\lambda}{2+\lambda})^2 \frac{R_{V/P}}{1+R_{V/P}} \langle N_M \rangle$	$\frac{3}{4} p_s p_{\bar{s}} \langle N_M \rangle$	$\frac{3}{4} (\frac{\lambda}{2+\lambda})^2 \langle N_M \rangle$
$\Lambda$	$\frac{6R_{O/D}}{1+2R_{O/D}} p_u^2 p_s \langle N_B \rangle$	$(\frac{5.30+12R_{O/D}}{1+2R_{O/D}} + \frac{5.64}{1+R_{O/D}}) p_u^2 p_s \langle N_B \rangle$	$\frac{\lambda}{(2+\lambda)^3} (\frac{5.30+12R_{O/D}}{1+2R_{O/D}} + \frac{5.64}{1+R_{O/D}}) \langle N_B \rangle$	$7.74 p_u^2 p_s \langle N_B \rangle$	$\frac{7.74\lambda}{(2+\lambda)^3} \langle N_B \rangle$
$\Sigma^+$	$\frac{3R_{O/D}}{1+R_{O/D}} p_u^2 p_s \langle N_B \rangle$	$(\frac{0.35}{1+2R_{O/D}} + \frac{0.18+3R_{O/D}}{1+R_{O/D}}) p_u^2 p_s \langle N_B \rangle$	$\frac{\lambda}{(2+\lambda)^3} (\frac{0.35}{1+2R_{O/D}} + \frac{0.18+3R_{O/D}}{1+R_{O/D}}) \langle N_B \rangle$	$2.13 p_u^2 p_s \langle N_B \rangle$	$\frac{2.13\lambda}{(2+\lambda)^3} \langle N_B \rangle$
$\Sigma^0$	$\frac{6R_{O/D}}{1+2R_{O/D}} p_u^2 p_s \langle N_B \rangle$	$(\frac{6R_{O/D}}{1+2R_{O/D}} + \frac{0.36}{1+R_{O/D}}) p_u^2 p_s \langle N_B \rangle$	$\frac{\lambda}{(2+\lambda)^3} (\frac{6R_{O/D}}{1+2R_{O/D}} + \frac{0.36}{1+R_{O/D}}) \langle N_B \rangle$	$2.52 p_u^2 p_s \langle N_B \rangle$	$\frac{2.52\lambda}{(2+\lambda)^3} \langle N_B \rangle$
$\Sigma^-$	$\frac{3R_{O/D}}{1+R_{O/D}} p_u^2 p_s \langle N_B \rangle$	$(\frac{0.35}{1+2R_{O/D}} + \frac{0.18+3R_{O/D}}{1+R_{O/D}}) p_u^2 p_s \langle N_B \rangle$	$\frac{\lambda}{(2+\lambda)^3} (\frac{0.35}{1+2R_{O/D}} + \frac{0.18+3R_{O/D}}{1+R_{O/D}}) \langle N_B \rangle$	$2.13 p_u^2 p_s \langle N_B \rangle$	$\frac{2.13\lambda}{(2+\lambda)^3} \langle N_B \rangle$
$\bar{\Lambda}$	$\frac{6R_{O/D}}{1+2R_{O/D}} p_{\bar{u}}^2 p_{\bar{s}} \langle N_{\bar{B}} \rangle$	$(\frac{5.30+12R_{O/D}}{1+2R_{O/D}} + \frac{5.64}{1+R_{O/D}}) p_{\bar{u}}^2 p_{\bar{s}} \langle N_{\bar{B}} \rangle$	$\frac{\lambda}{(2+\lambda)^3} (\frac{5.30+12R_{O/D}}{1+2R_{O/D}} + \frac{5.64}{1+R_{O/D}}) \langle N_B \rangle$	$7.74 p_{\bar{u}}^2 p_{\bar{s}} \langle N_{\bar{B}} \rangle$	$\frac{7.74\lambda}{(2+\lambda)^3} \langle N_B \rangle$
$\bar{\Sigma}^-$	$\frac{3R_{O/D}}{1+R_{O/D}} p_{\bar{u}}^2 p_{\bar{s}} \langle N_{\bar{B}} \rangle$	$(\frac{0.35}{1+2R_{O/D}} + \frac{0.18+3R_{O/D}}{1+R_{O/D}}) p_{\bar{u}}^2 p_{\bar{s}} \langle N_{\bar{B}} \rangle$	$\frac{\lambda}{(2+\lambda)^3} (\frac{0.35}{1+2R_{O/D}} + \frac{0.18+3R_{O/D}}{1+R_{O/D}}) \langle N_B \rangle$	$2.13 p_{\bar{u}}^2 p_{\bar{s}} \langle N_{\bar{B}} \rangle$	$\frac{2.13\lambda}{(2+\lambda)^3} \langle N_B \rangle$
$\bar{\Sigma}^0$	$\frac{6R_{O/D}}{1+2R_{O/D}} p_{\bar{u}}^2 p_{\bar{s}} \langle N_{\bar{B}} \rangle$	$(\frac{6R_{O/D}}{1+2R_{O/D}} + \frac{0.36}{1+R_{O/D}}) p_{\bar{u}}^2 p_{\bar{s}} \langle N_{\bar{B}} \rangle$	$\frac{\lambda}{(2+\lambda)^3} (\frac{6R_{O/D}}{1+2R_{O/D}} + \frac{0.36}{1+R_{O/D}}) \langle N_B \rangle$	$2.52 p_{\bar{u}}^2 p_{\bar{s}} \langle N_{\bar{B}} \rangle$	$\frac{2.52\lambda}{(2+\lambda)^3} \langle N_B \rangle$
$\bar{\Sigma}^+$	$\frac{3R_{O/D}}{1+R_{O/D}} p_{\bar{u}}^2 p_{\bar{s}} \langle N_{\bar{B}} \rangle$	$(\frac{0.35}{1+2R_{O/D}} + \frac{0.18+3R_{O/D}}{1+R_{O/D}}) p_{\bar{u}}^2 p_{\bar{s}} \langle N_{\bar{B}} \rangle$	$\frac{\lambda}{(2+\lambda)^3} (\frac{0.35}{1+2R_{O/D}} + \frac{0.18+3R_{O/D}}{1+R_{O/D}}) \langle N_B \rangle$	$2.13 p_{\bar{u}}^2 p_{\bar{s}} \langle N_{\bar{B}} \rangle$	$\frac{2.13\lambda}{(2+\lambda)^3} \langle N_B \rangle$

We do not list the corresponding results for pions in Table I since the corresponding expressions are quite long. This is because the pion receives contributions from the

decays of almost all the other mesons, baryons, and antibaryons. For example, for  $\pi^+$ , we have

$$\begin{aligned}
\langle N_{\pi^+}^f \rangle = & \langle N_{\pi^+} \rangle + \langle N_{\rho^+} \rangle + \langle N_{\rho^0} \rangle + \frac{2}{3} (\langle N_{\bar{K}^{*0}} \rangle + \langle N_{K^{*+}} \rangle) + 0.2734 \langle N_{\eta} \rangle + 0.9073 \langle N_{\omega} \rangle + 0.9274 \langle N_{\eta'} \rangle + 0.1568 \langle N_{\phi} \rangle \\
& + 0.94 (\langle N_{\Sigma^{*+}} \rangle + \langle N_{\bar{\Sigma}^{*+}} \rangle) + 0.0585 (\langle N_{\Sigma^{*0}} \rangle + \langle N_{\bar{\Sigma}^{*0}} \rangle) + \frac{2}{3} (\langle N_{\Xi^{*0}} \rangle + \langle N_{\bar{\Xi}^{*+}} \rangle) + \langle N_{\Delta^{*+}} \rangle + \frac{1}{3} (\langle N_{\Delta^+} \rangle + \langle N_{\bar{\Delta}^0} \rangle) + \langle N_{\bar{\Delta}^+} \rangle. \quad (53)
\end{aligned}$$

By inserting the results given by Eqs. (42) and (43), we obtain

$$\begin{aligned} \langle N_{\pi^+}^f \rangle = & \frac{1.71 + 2.91R_{V/P}}{1 + R_{V/P}} p_u p_{\bar{u}} \langle N_M \rangle + \frac{0.49 + 0.16R_{V/P}}{1 + R_{V/P}} p_s p_{\bar{s}} \langle N_M \rangle + \frac{2}{3} \frac{R_{V/P}}{1 + R_{V/P}} p_{\bar{u}} p_s \langle N_M \rangle \\ & + \frac{2}{3} \frac{R_{V/P}}{1 + R_{V/P}} p_u p_{\bar{s}} \langle N_M \rangle + \left( \frac{2.82}{1 + R_{O/D}} + \frac{0.35}{1 + 2R_{O/D}} \right) p_u^2 p_s \langle N_B \rangle + \left( \frac{2.82}{1 + R_{O/D}} + \frac{0.35}{1 + 2R_{O/D}} \right) p_{\bar{u}}^2 p_{\bar{s}} \langle N_{\bar{B}} \rangle \\ & + \frac{2}{1 + R_{O/D}} p_u p_s^2 \langle N_B \rangle + \frac{2}{1 + R_{O/D}} p_{\bar{u}} p_{\bar{s}}^2 \langle N_{\bar{B}} \rangle + \frac{2 + R_{O/D}}{1 + R_{O/D}} p_u^3 \langle N_B \rangle + \frac{2 + R_{O/D}}{1 + R_{O/D}} p_{\bar{u}}^3 \langle N_{\bar{B}} \rangle. \end{aligned} \quad (54)$$

For systems without net quarks, we have  $p_u = p_{\bar{u}} = p_d = p_{\bar{d}}$  and  $p_s = p_{\bar{s}}$ , and thus we obtain

$$\begin{aligned} \langle N_{\pi^+}^f \rangle = & \frac{1.71 + 2.91R_{V/P}}{1 + R_{V/P}} p_u^2 \langle N_M \rangle + \frac{0.49 + 0.16R_{V/P}}{1 + R_{V/P}} p_s^2 \langle N_M \rangle + \frac{4}{3} \frac{R_{V/P}}{1 + R_{V/P}} p_u p_s \langle N_M \rangle \\ & + \left( \frac{5.64}{1 + R_{O/D}} + \frac{0.70}{1 + 2R_{O/D}} \right) p_u^2 p_s \langle N_B \rangle + \frac{4}{1 + R_{O/D}} p_u p_s^2 \langle N_B \rangle + \frac{4 + 2R_{O/D}}{1 + R_{O/D}} p_u^3 \langle N_B \rangle. \end{aligned} \quad (55)$$

If we take  $R_{V/P} = 3$  and  $R_{O/D} = 2$ , we have

$$\begin{aligned} \langle N_{\pi^+}^f \rangle = & 2.61 p_u^2 \langle N_M \rangle + p_u p_s \langle N_M \rangle + 0.24 p_s^2 \langle N_M \rangle + \frac{4}{3} p_u p_s^2 \langle N_B \rangle + 2.02 p_u^2 p_s \langle N_B \rangle + \frac{8}{3} p_u^3 \langle N_B \rangle \\ = & \frac{2.61 + \lambda + 0.24\lambda^2}{(2 + \lambda)^2} \langle N_M \rangle + \frac{8/3 + 2.02\lambda + 4/3\lambda^2}{(2 + \lambda)^3} \langle N_B \rangle. \end{aligned} \quad (56)$$

From these results, we see clearly that there exist many simple relationships between the yields of different hadrons. These are the characteristics for hadron production in the combination mechanism. We will list some of these simple relations in the following. Before doing that, we first discuss the net quark influences in the next section.

### E. Influence of the net quarks

In a heavy-ion collision at high energy, the produced quark-antiquark system consisting of the newly produced quarks and antiquarks and the net quarks from the incident nuclei. For a subsample of this quark-antiquark system in a given kinematic region, we have, in general,

$$\langle N_q \rangle = \langle N_{\bar{q}} \rangle + \langle N_q^{net} \rangle. \quad (57)$$

Both the momentum and flavor distributions of these net quarks are different from those for the newly produced ones, and this leads to observable effects in the final hadrons produced in hadronization. We expect that they have influences on the following aspects:

(i) The difference in momentum distribution leads to different  $\gamma_M(N_q, N_{\bar{q}}, \sqrt{s})$  and  $\gamma_B(N_q, N_{\bar{q}}, \sqrt{s})$ , as seen clearly from Eqs. (22) and (23). Furthermore, since  $\langle N_q \rangle > \langle N_{\bar{q}} \rangle$ , the average number of baryons,  $\langle N_B \rangle$ , should be accordingly larger than  $\langle N_{\bar{B}} \rangle$ . The ratio  $\langle N_B \rangle / \langle N_M \rangle$  should in general depend on  $\langle N_q^{net} \rangle / \langle N_q \rangle$ .

(ii) The distribution of the number of quarks,  $N_q$ , at a given  $\langle N_q \rangle$  is different from the corresponding distribution of the antiquarks. The distribution of the number of the net quarks,  $N_q^{net}$ , at a given  $\langle N_q^{net} \rangle$  is different from those

for the newly produced quarks and/or antiquarks. This leads to a difference between the quark number distribution  $P_q(N_q, \langle N_q \rangle, \sqrt{s})$  and the antiquark number distribution  $P_{\bar{q}}(N_{\bar{q}}, \langle N_{\bar{q}} \rangle, \sqrt{s})$ .

(iii) The flavor distribution of quarks is different from that for the antiquarks. The net quarks take only two flavors,  $u$  and  $d$ . At a given  $N_q^{net}$ , the average numbers of  $u$  and  $d$  net quarks are determined by the numbers of protons and neutrons in the incident nuclei. For a given AA collision,  $\bar{N}_u^{net} : \bar{N}_d^{net} = (A + Z) : (2A - Z)$ , where  $A$  and  $Z$  are the numbers of nucleons and protons, respectively, in the incident nucleus  $A$ . The numbers  $N_u^{net}$  and  $N_d^{net}$  follow a binominal distribution with  $p_u^{net} : p_d^{net} = (A + Z) : (2A - Z)$ . For the newly produced quarks or antiquarks, at a given  $N_q^{new}$  or  $N_{\bar{q}}$ , the numbers of them of different flavors follow the multinominal distribution as given by Eq. (35). Hence, including the net quark contribution, the distribution of the numbers ( $N_u$ ,  $N_d$ , and  $N_s$ ) of the different flavors ( $u$ ,  $d$ , and  $s$ ) of quarks at a given number of quarks ( $N_q = N_u + N_d + N_s$ ) is also different from the corresponding distribution for the antiquarks.

The detailed calculations of the influences of these effects on the hadron yield ratios in the combination mechanism depend on the particular models. In this paper, we present a rough estimate of, at least, the qualitative tendency of these effects by using the following two approximations.

First, we approximate that the flavor distribution of the number of quarks in the subsample of the system is independent of that for the antiquarks. That for the antiquarks is given by the multinominal given by Eq. (36). For the quarks, we approximate it by a multinominal distribution

in the same form as that for the antiquarks but with different probabilities of the different flavors. For the newly produced quarks, the flavor distribution should be  $1 : 1 : \lambda$  for  $u, d$ , and  $s$  quarks, the same as those for the antiquarks, but for the net quarks, it should be  $(A + Z) : (2A - Z) : 0$ . We take the average and obtain, at given  $N_q$  and  $N_q^{net}$ ,

$$p_u = \frac{1}{2 + \lambda} \left(1 - \frac{N_q^{net}}{N_q}\right) + \frac{A + Z}{3A} \frac{N_q^{net}}{N_q}, \quad (58)$$

$$p_d = \frac{1}{2 + \lambda} \left(1 - \frac{N_q^{net}}{N_q}\right) + \frac{2A - Z}{3A} \frac{N_q^{net}}{N_q}, \quad (59)$$

$$p_s = \frac{\lambda}{2 + \lambda} \left(1 - \frac{N_q^{net}}{N_q}\right). \quad (60)$$

In the case in which the  $u$  and  $d$  difference is not very large, we neglect it and consider the case in which  $\overline{N}_u^{net} : \overline{N}_d^{net} = 1 : 1$ . In this case, we have, at given  $N_q$  and  $N_q^{net}$ ,  $p_u : p_d : p_s = 1 : 1 : \lambda_q$  and

$$\lambda_q = \frac{\overline{N}_s}{\overline{N}_u} = \lambda \left[1 + \left(1 + \frac{\lambda}{2}\right) \frac{N_q^{net}}{N_q - N_q^{net}}\right]^{-1}. \quad (61)$$

Under this approximation, we see that Eqs. (40) and (41), the results we obtained in Sec. IIC for directly produced hadrons in the combination of a quark-antiquark system at given  $N_q$  and  $N_{\bar{q}}$ , are still valid. We only need to note that the  $p_{q_i}$  in this case is a function of  $N_q$  and  $N_q^{net}$  as given by Eq. (61). The influence of the isospin violation in net quarks can manifest itself in the difference between the average yields of hadrons belonging to the same charge multiplet. This can be studied separately in experiments. Since our purpose is a rough estimation of the net quark influence, we consider in the following first the simplified case where  $u$  and  $d$  are equal but leave the isospin difference for future studies.

Second, for a subsample in a given kinematic region of the bulk system produced in a given AA collision at given energy  $\sqrt{s}$ , the averages  $\langle N_q \rangle$  and  $\langle N_{\bar{q}} \rangle$  are fixed and the numbers  $N_q$  and  $N_{\bar{q}}$  follow the distribution  $P(N_q, N_{\bar{q}}; \langle N_q \rangle, \langle N_{\bar{q}} \rangle, \sqrt{s})$ . The average numbers of the hadrons produced should be the average over this distribution. In general such averages depend on the precise form of  $P(N_q, N_{\bar{q}}; \langle N_q \rangle, \langle N_{\bar{q}} \rangle, \sqrt{s})$ . In the rough estimations we made here, we approximate these averages by taking the corresponding values of the quantities at the averages  $\langle N_q \rangle$  and  $\langle N_{\bar{q}} \rangle$ , i.e.,

$$\langle N_{h_j}(N_q, N_{\bar{q}}, \sqrt{s}) \rangle \approx \overline{N}_{h_j}(\langle N_q \rangle, \langle N_{\bar{q}} \rangle, \sqrt{s}). \quad (62)$$

Under these two approximations, all the results presented in the last four sections where we distinguish between  $p_{q_i}$  and  $p_{\bar{q}_i}$  apply and we can use them to make estimates of the effects of the net quarks.

## F. Ratios of the yields of different hadrons

From the results given in Table I, we see that there are many simple relations between the yields of different

hadrons. In particular, for the directly produced hadrons, such relationships are very simple. Even for the final hadrons, although the decay influences are often very large, there still exists a set of simple relations between them. For example, independent of the values of  $R_{V/P}$  and  $R_{O/D}$ , for the final hadrons where contributions from strong and electromagnetic decays are taken into account, we have

$$\frac{\langle N_{\Xi^-}^f \rangle}{\langle N_p^f \rangle} = \frac{\langle N_{\Xi^0}^f \rangle}{\langle N_p^f \rangle} = \frac{3}{4} \lambda_q^2, \quad (63)$$

$$\frac{\langle N_{\Omega^-}^f \rangle}{\langle N_p^f \rangle} = \frac{1}{4} \lambda_q^3, \quad (64)$$

$$\frac{\langle N_{\Xi^+}^f \rangle}{\langle N_{\bar{p}}^f \rangle} = \frac{\langle N_{\Xi^0}^f \rangle}{\langle N_{\bar{p}}^f \rangle} = \frac{3}{4} \lambda^2, \quad (65)$$

$$\frac{\langle N_{\Omega^+}^f \rangle}{\langle N_{\bar{p}}^f \rangle} = \frac{1}{4} \lambda^3, \quad (66)$$

$$\frac{\langle N_{\bar{p}}^f \rangle}{\langle N_p^f \rangle} = \left(\frac{2 + \lambda_q}{2 + \lambda}\right)^3 \frac{\langle N_{\bar{B}} \rangle}{\langle N_B \rangle}, \quad (67)$$

$$\frac{\langle N_{\bar{\Lambda}}^f \rangle}{\langle N_{\Lambda}^f \rangle} = \left(\frac{2 + \lambda_q}{2 + \lambda}\right)^3 \frac{\lambda}{\lambda_q} \frac{\langle N_{\bar{B}} \rangle}{\langle N_B \rangle}, \quad (68)$$

$$\frac{\langle N_{\Xi^+}^f \rangle}{\langle N_{\Xi^-}^f \rangle} = \left(\frac{2 + \lambda_q}{2 + \lambda}\right)^3 \left(\frac{\lambda}{\lambda_q}\right)^2 \frac{\langle N_{\bar{B}} \rangle}{\langle N_B \rangle}, \quad (69)$$

$$\frac{\langle N_{\Omega^+}^f \rangle}{\langle N_{\Omega^-}^f \rangle} = \left(\frac{2 + \lambda_q}{2 + \lambda}\right)^3 \left(\frac{\lambda}{\lambda_q}\right)^3 \frac{\langle N_{\bar{B}} \rangle}{\langle N_B \rangle}. \quad (70)$$

In the case in which net quark contribution is negligible, we have,  $\lambda_q = \lambda$  and  $\langle N_B \rangle = \langle N_{\bar{B}} \rangle$ , so that

$$\frac{\langle N_{\Xi^-}^f \rangle}{\langle N_p^f \rangle} = \frac{\langle N_{\Xi^0}^f \rangle}{\langle N_p^f \rangle} = \frac{3}{4} \lambda^2, \quad (71)$$

$$\frac{\langle N_{\Omega^-}^f \rangle}{\langle N_p^f \rangle} = \frac{\langle N_{\Omega^+}^f \rangle}{\langle N_{\bar{p}}^f \rangle} = \frac{1}{4} \lambda^3, \quad (72)$$

and  $\langle N_{\bar{B}_j}^f \rangle / \langle N_{B_j}^f \rangle = 1$  for all the different types of  $B_j$ .

In the case in which  $R_{V/P} = 3$  and  $R_{O/D} = 2$ , we have more such simple relations such as

$$\frac{\langle N_{K^-}^f \rangle}{\langle N_{K^+}^f \rangle} = \left(\frac{\lambda_q}{\lambda}\right) \left(\frac{1 + 0.37\lambda}{1 + 0.37\lambda_q}\right), \quad (73)$$

$$\frac{\langle N_{\Lambda}^f \rangle}{\langle N_{\bar{p}}^f \rangle} = 1.935\lambda, \quad (74)$$

$$\frac{\langle N_{\Lambda}^f \rangle}{\langle N_p^f \rangle} = 1.935\lambda_q. \quad (75)$$



These relations are intrinsic properties of the combination models in the sense that they do not depend on the details of particular combination models but are determined mainly by the basic ideas of the combination mechanism. They can be used to determine the free parameters and/or to test the mechanism. We note in particular the following two features.

(i) For a system with equal average numbers of quarks and antiquarks, i.e., where net quark contributions are negligible, these ratios are quite simple and can be tested by the data in extremely high energy collisions, e.g., at the LHC.

(ii) With net quark contributions, most of these yield ratios of different hadrons depend on  $\langle N_q^{net} \rangle / \langle N_q \rangle$ . We see in particular that the ratios of the yields of hadrons to those of the corresponding antihadrons deviate from one in general. They tend to one for reactions at very high energies where the net quark contribution tends to vanish. This leads to an energy dependence of such ratios, even for particles such as  $\Omega^-$  and  $\bar{\Omega}^+$ . Such a property for the combination mechanism is different from what one expects from fragmentation and can be used as a check to differentiate the different hadronization mechanisms.

We can also build some combinations of the average yields of the hadrons and obtain simple results for some more sophisticated ratios such as,

$$A \equiv \frac{\langle N_{\Lambda} \rangle \langle N_{K^-} \rangle \langle N_p \rangle}{\langle N_{\Lambda} \rangle \langle N_{K^+} \rangle \langle N_{\bar{p}} \rangle} = 1, \quad (76)$$

$$B \equiv \frac{\langle N_{\Lambda} \rangle \langle N_{K^-} \rangle \langle N_{\Xi^+} \rangle}{\langle N_{\Lambda} \rangle \langle N_{K^+} \rangle \langle N_{\Xi^-} \rangle} = 1, \quad (77)$$

$$d_{\Xi}^{\Lambda} \equiv \frac{\langle N_{\Lambda} \rangle \langle N_{\Xi^+} \rangle}{\langle N_{\Lambda} \rangle \langle N_{\Xi^-} \rangle} = \left( \frac{2 + \lambda_q}{2 + \lambda} \right)^6 \left( \frac{\lambda}{\lambda_q} \right)^3 \left( \frac{\langle N_B \rangle}{\langle N_{\bar{B}} \rangle} \right)^2, \quad (78)$$

$$d_{\Omega}^p \equiv \frac{\langle N_{\bar{p}} \rangle \langle N_{\bar{\Omega}^+} \rangle}{\langle N_p \rangle \langle N_{\Omega^-} \rangle} = d_{\Xi}^{\Lambda}. \quad (79)$$

We see that they all lead to simple results and these relations are not influenced by the resonance decays except  $\phi \rightarrow K^+ K^-$ , which slightly changes  $A$  and  $B$  from unity. For the case in which  $N_q^{net} = 0$ , all four ratios are equal to unity.

As a brief summary, we emphasize once more that the model that we consider in this section is intended to be a general case based on the basic ideas of the combination mechanism. The purpose is to concentrate on the hadron yield correlations in the combination models. No effort is made to study other properties such as momentum distribution, etc. The results obtained follow from the basic ideas of the combination mechanism and a number of assumptions, simplifications, and/or approximations such as the factorization of flavor and momentum dependence of the kernel functions, the flavor independence of the quark-antiquark momentum distribution, the independent production of different flavors of quarks and antiquarks, and the approximations made in considering the net quark influences. These results do not depend on the detailed form of the momentum dependence of kernel functions

and/or the momentum distributions of the quarks and antiquarks. They even do not depend on whether the quark number conservation or depletion is imposed in the combination process. Such a conservation or depletion of quark number influences the relationship between the average number of mesons (or baryons, or antibaryons) produced and the number of quarks and/or antiquarks participating in the combination process (see in particular the discussion in [40]) but does not influence the hadron yield ratios if the factorization is assumed. These results should also be valid in the combination models discussed in the literature wherever these assumptions and/or approximations are also made (explicitly or implicitly). Some of them should even be common in these different models [9–13, 23]. In fact, some of the relationships presented above have also been derived in these literature [9–13, 23]. For example, relations similar to those given in Eqs. (76)–(79) have been obtained in [9]. These relations can be used to test the validity of the combination mechanism and the assumptions made. We will compare them with the data available in the next section.

### III. COMPARISON WITH DATA

There are already quite abundant data available from experiments in a quite broad energy region, from low SPS energies to RHIC and LHC energies [18, 21, 22, 24–36]. We compare the results obtained in the last section with these data in the following.

As mentioned earlier, the study presented in Sec. II is intended to be a general case for the combination of a quark-antiquark system consisting of  $N_q$  quarks and  $N_{\bar{q}}$  antiquarks. No effort is made to ascertain whether the combination mechanism dominates the production of hadrons in the given kinematic region in AA collisions. In this section, we choose the data in the central rapidity regions in different AA collisions at different energies. The agreement and/or disagreement of the theoretical results with these data should give us a signature of whether the combination mechanism with the above-mentioned assumptions and/or approximations is applicable. Also, since some of the theoretical results depend on more inputs and some of them depend on fewer inputs, we make the comparison at different levels.

#### A. Comparison with LHC data

At the first level, we consider a subsample of the quark-antiquark system in the central rapidity region produced in AA collisions at very high energies. We suppose the energies are very high and the subsample that we consider is only a small part of the whole quark-antiquark system produced in the collision process so that the influence of the net quarks and that from the global flavor compensation are negligible. In this limiting case, the results for the ratios of the yields of different hadrons are divided into three

classes.

In the first class, we consider the ratios of the yields of hadrons to those of the corresponding antihadrons. Such ratios are unity, independent of any parameter. This can be considered as a criterion for the validity of this limiting case. Results from LHC experiments can be considered as an example for this case. In the first three lines of Table II, we show the available experimental results for the particle to antiparticle ratios such as  $\langle N_{\pi^-}^f \rangle / \langle N_{\pi^+}^f \rangle$ ,  $\langle N_{K^-}^f \rangle / \langle N_{K^+}^f \rangle$ , and  $\langle N_{\bar{p}}^f \rangle / \langle N_p^f \rangle$  at mid-rapidity. The data are obtained from Ref. [21]. We see that they are indeed very close to unity.

TABLE II: Hadron yield ratios obtained for  $N_q^{\text{net}} = 0$  compared with data from the LHC in Pb + Pb collisions at  $\sqrt{s} = 2.76$  TeV. The data are taken from Refs. [21, 22]. The experimental result for  $K^-/\pi^-$  is used to determine the strangeness suppression factor  $\lambda$ .

Ratios	Data	Calculations
$\pi^-/\pi^+$	$1.000 \pm 0.080$	1
$K^-/K^+$	$0.987 \pm 0.076$	1
$\bar{p}/p$	$0.995 \pm 0.077$	1
$\phi/K^+$	—	0.278
$K_S^0/K^+$	—	0.959
$\Lambda/p$	—	0.832
$\Xi^-/p$	—	0.139
$\Omega^-/p$	—	0.020
$K^-/\pi^-$	$0.155 \pm 0.012$	0.155
$p/\pi^+$	$0.045 \pm 0.004$	0.043
$\Lambda/\pi^+$	—	0.036
$\Xi^-/\pi^+$	$0.005 \pm 0.001$	0.006
$\Omega^-/\pi^+$	$0.001 \pm 0.0002$	0.001
$p/K^+$	—	0.275
$\Lambda/K^+$	—	0.229
$\Xi^-/K^+$	—	0.038
$\Omega^-/K^+$	—	0.005

In the second class, we consider the ratios of the average yields of hadrons such as  $\phi/K^+$ ,  $\Lambda/p$ , and so on as shown in the second part of Table II (from the forth line to the eighth line). Because in this case we have  $\langle N_B \rangle = \langle N_{\bar{B}} \rangle$  and  $\lambda_q = \lambda$ , these ratios depend only on one free parameter, the strangeness suppression factor  $\lambda$ , and are independent of the particular models. We can fix  $\lambda$  by using the data for one ratio and make predictions for other particle ratios. Such results can be used to check the validity of the combination picture.

In the third class, we consider the ratio of a specified meson to a specified baryon. To obtain the results for such ratios, we need the input for  $\langle N_B \rangle / \langle N_M \rangle$ . This can be slightly different in different combination models. As an example, in the third part of Table II (from the ninth line to the end), we show the results obtained by taking

$$\langle N_B \rangle / \langle N_M \rangle = 1/12. \quad (80)$$

This is obtained by parametrizing the results for large

$N_q = N_{\bar{q}}$  of the Monte Carlo generator (SDQCM) based on the combination rule developed by the Shandong group [13, 15, 37] which has reproduced the data well. This parametrization is valid with high accuracy for  $N_q$  larger than, say, 100. The parameter  $\lambda$  is taken as  $\lambda = 0.43$  by fitting the data of  $K^-/\pi^-$ . We see that the theoretical results agree with the LHC data whenever available.

## B. Comparison with RHIC and SPS data

At the second level of comparison, we consider the case where  $N_q^{\text{net}} \neq 0$ . In this case, we need a further input of  $\langle N_q^{\text{net}} \rangle / \langle N_q \rangle$  which describes the strength of the net quark influence. Clearly, this ratio depends on the type and the energy of the incident nuclei, and also on the kinematic region that we consider. For example, to make a good comparison with the data from the RHIC and the SPS [18, 24–36], we need to take such effects into account.

In practice, to carry out the calculations, we first fix  $\lambda_q$  and  $\lambda$  by using the data for  $\langle N_{\Xi^-} \rangle / \langle N_p \rangle$  to determine  $\lambda_q$  and  $\langle N_{\Xi^+} \rangle / \langle N_{\bar{p}} \rangle$  to determine  $\lambda$ , and we then derive  $\langle N_q^{\text{net}} \rangle / \langle N_q \rangle$  using Eq. (61). In Table III, we show the results obtained by taking the data for  $dN/dy$  at  $y = 0$  (where  $y$  denotes the rapidity of the hadron). These results can be used to calculate the hadron ratios in the mid-rapidity regions.

TABLE III: The fixed  $\lambda$  and  $\lambda_q$  and the derived  $\langle N_q^{\text{net}} \rangle / \langle N_q \rangle$  at RHIC and SPS energies.

Energy	RHIC $\sqrt{s}$ (GeV)			SPS $E_{\text{beam}}$ (A GeV)			
	200	130	62.4	158	80	40	30
$\lambda$	0.425	0.412	0.369	0.499	0.567	0.540	0.645
$\lambda_q$	0.397	0.375	0.328	0.255	0.232	0.193	0.192
$\langle N_q^{\text{net}} \rangle / \langle N_q \rangle$	0.056	0.074	0.095	0.434	0.529	0.587	0.640

From the table, we see that  $\langle N_q^{\text{net}} \rangle / \langle N_q \rangle$  is small at RHIC energies but quite large at SPS energies and the expected effects should be large at those energies.

Using the obtained  $\lambda$  and  $\lambda_q$  values as inputs, we calculate the hadron yield ratios that are independent of  $\langle N_B \rangle / \langle N_M \rangle$ . The results are given in Table IV. The corresponding data are from Refs. [18, 24–35].

We see in particular that the ratio of the yield of hadrons to that of the corresponding antihadrons is not unity in this case. We also note that the ratio such as  $\langle N_{K^-}^f \rangle / \langle N_{K^+}^f \rangle = (\lambda_q/\lambda)(1 + 0.37\lambda)/(1 + 0.37\lambda_q)$  is a good example to show the change of the effective strange suppression for quarks. The results should decrease monotonically with increasing  $\langle N_q^{\text{net}} \rangle / \langle N_q \rangle$ . There are data available for  $\langle N_{K^-}^f \rangle / \langle N_{K^+}^f \rangle$  at different energies [18, 24, 25] and the data show clearly that the ratio increases with increasing energy. This qualitative tendency is consistent with the effect of net quark contribution since the relative influence of the net quarks becomes smaller at higher energies.

To calculate other ratios, we need  $\langle N_B \rangle / \langle N_M \rangle$  and

TABLE IV: Comparison of the calculated results for hadron yield ratios that are independent of  $\langle N_B \rangle / \langle N_M \rangle$  with data from RHIC and SPS energies [18, 24–35]. The experimental result for  $\Xi^-/p$  and that for  $\bar{\Xi}^+/\bar{p}$  are used to determine the strangeness suppression factor  $\lambda_q$  and  $\lambda$ , respectively.

Ratio	$\sqrt{s}$ (GeV) for Au + Au at RHIC						$E_{beam}$ (A GeV) for Pb + Pb at SPS							
	200		130		62.4		158		80		40		30	
	data	theory	data	theory	data	theory	data	theory	data	theory	data	theory	data	theory
$\bar{\Lambda}/\bar{p}$	$0.94 \pm 0.15$	0.82	$0.93 \pm 0.34$	0.80	$0.82 \pm 0.18$	0.71	$0.98 \pm 0.23$	0.97	$1.22 \pm 0.30$	1.10	$1.31 \pm 0.35$	1.05	$1.31 \pm 0.41$	1.25
$\bar{\Xi}^+/\bar{p}$	$0.14 \pm 0.03$	0.14	$0.13 \pm 0.03$	0.13	$0.10 \pm 0.03$	0.10	$0.19 \pm 0.05$	0.19	$0.24 \pm 0.07$	0.24	$0.22 \pm 0.07$	0.22	$0.31 \pm 0.15$	0.31
$\bar{\Omega}^+/\bar{p}$	—	0.019	—	0.017	$0.017 \pm 0.005$	0.013	$0.042 \pm 0.020$	0.031	—	0.046	—	0.039	—	0.067
$\Lambda/p$	$0.91 \pm 0.15$	0.77	$0.90 \pm 0.30$	0.73	$0.78 \pm 0.16$	0.63	$0.37 \pm 0.09$	0.49	$0.45 \pm 0.08$	0.45	$0.37 \pm 0.06$	0.37	$0.35 \pm 0.06$	0.37
$\Xi^-/p$	$0.12 \pm 0.02$	0.12	$0.11 \pm 0.03$	0.11	$0.08 \pm 0.02$	0.08	$0.05 \pm 0.01$	0.05	$0.04 \pm 0.01$	0.04	$0.03 \pm 0.01$	0.03	$0.03 \pm 0.01$	0.03
$\Omega^-/p$	—	0.016	—	0.013	$0.010 \pm 0.003$	0.009	$0.005 \pm 0.001$	0.004	—	0.003	—	0.002	—	0.002
$K^-/K^+$	$0.96 \pm 0.17$	0.94	$0.92 \pm 0.09$	0.92	$0.86 \pm 0.09$	0.90	$0.57 \pm 0.05$	0.55	$0.48 \pm 0.04$	0.46	$0.38 \pm 0.04$	0.40	$0.37 \pm 0.04$	0.34
$\bar{K}^0/K^0$	—	0.94	—	0.92	—	0.90	—	0.54	—	0.44	—	0.39	—	0.33
$K^0/K^+$	—	0.96	—	0.96	—	0.97	—	0.97	—	0.98	—	0.98	—	0.98

$\langle N_{\bar{B}} \rangle / \langle N_M \rangle$ . For this purpose, we parametrize  $\langle N_{\bar{B}} \rangle / \langle N_M \rangle$  for the case where  $N_q^{net} \neq 0$  using SDQCM [13, 15, 37], and we obtain

$$\frac{\langle N_{\bar{B}} \rangle}{\langle N_M \rangle} = \frac{1}{12} \left( 1 - \frac{\langle N_q^{net} \rangle}{\langle N_q \rangle} \right)^{2.8}. \quad (81)$$

We found that this parametrization is a good approximation to the results obtained from SDQCM for different  $N_q$  in the range of  $N_q > 100$ . In this model, quark number conservation or depletion is imposed in the combination process so that  $3\langle N_B \rangle + \langle N_M \rangle = \langle N_q \rangle$  and  $3\langle N_{\bar{B}} \rangle + \langle N_M \rangle = \langle N_q \rangle - \langle N_q^{net} \rangle$ , and we derive

$$\frac{\langle N_{\bar{B}} \rangle}{\langle N_M \rangle} = \left( \frac{\langle N_{\bar{B}} \rangle}{\langle N_M \rangle} + \frac{1}{3} \frac{\langle N_q^{net} \rangle}{\langle N_q \rangle} \right) / \left( 1 - \frac{\langle N_q^{net} \rangle}{\langle N_q \rangle} \right). \quad (82)$$

With Eqs. (81) and (82), we calculate the ratios of the yields of different baryons to mesons. The results are shown in Table V. The data are calculated from  $dN/dy$  at  $y = 0$  for different energies from Refs. [18, 24–35].

With the values of  $\lambda$  and  $\lambda_q$  given in Table III and  $\langle N_{\bar{B}} \rangle / \langle N_B \rangle$  derived from Eqs. (81) and (82), we calculate in particular the ratios of the antibaryons to the corresponding baryons at different energies. The results obtained are shown in Fig. 1 with the open symbols connected by different lines to guide the eye. The filled symbols with error bars are the experimental data taken from Refs. [24, 26–28, 36]. From Fig. 1, we see in particular that  $\langle N_{\bar{\Omega}^+}^f \rangle / \langle N_{\bar{\Xi}^+}^f \rangle > \langle N_{\bar{\Xi}^+}^f \rangle / \langle N_{\bar{\Lambda}^+}^f \rangle > \langle N_{\bar{\Lambda}^+}^f \rangle / \langle N_{\bar{p}}^f \rangle > \langle N_{\bar{p}}^f \rangle / \langle N_p^f \rangle$ . This was first pointed out by the NA49 Collaboration and was regarded as a distinct hierarchy of the antibaryon to baryon ratios [27]. The hierarchy can be naturally reproduced by Eqs. (67)–(70) in the simple quark combination

models. We also see that the antibaryon to baryon ratios increase with increasing energy, indicating that the net quark influence becomes smaller at higher energies.

It was considered as a surprise that the net quark also influences  $\langle N_{\bar{\Omega}^+}^f \rangle / \langle N_{\bar{\Xi}^+}^f \rangle$  significantly, as shown in Fig. 1, although the  $\bar{\Omega}$  hyperon does not consist of  $u$  or  $d$  quarks. We see that from the lowest SPS to the highest RHIC energy,  $\langle N_{\bar{\Omega}^+}^f \rangle / \langle N_{\bar{\Xi}^+}^f \rangle$  increases from about 0.4 to unity. This cannot be understood in the fragmentation models but can be naturally explained in the framework of quark combination models. Because in the combination models more net quarks imply more chances for the antistrange quarks to meet quarks to form mesons, there will be more antistrange quarks exhausted to form kaons than strange quarks. The probability for a strange antiquark to combine with other antiquarks to form an antibaryon is smaller than that for a strange quark to combine with other quarks to form a baryon. This leads to fewer  $\bar{\Omega}$ 's than  $\Omega$ 's for the same number of strange quarks and antistrange quarks. This qualitative tendency is consistent with data, and from the figure we see also the quantitative results agree well with the data.

We also compare the experimental results with the predictions shown in Eqs. (76)–(79) for the “more sophisticated” ratios  $A = B = 1$  and  $d_{\Xi}^{\Lambda} = d_{\Omega}^p$ . From the data available [24, 26–28, 36], we calculate these ratios and show the results in Table VI and Fig. 2 for  $d_{\Xi}^{\Lambda}$  and  $d_{\Omega}^p$  and  $A$  and  $B$ , respectively. We see that the data are consistent with  $d_{\Xi}^{\Lambda} = d_{\Omega}^p$ . The results for  $A$  and  $B$  at RHIC energies are consistent with unity while the error bars for those at SPS energies are too large to make a judgment.

TABLE V: Comparison of the calculated results for the baryon-to-meson ratios that are dependent of  $\langle N_B \rangle / \langle N_M \rangle$  with data from RHIC and SPS energies [18, 24–35].

Ratio	$\sqrt{s}$ (GeV) for Au + Au at RHIC						$E_{beam}$ (A GeV) for Pb + Pb at SPS							
	200		130		62.4		158		80		40		30	
	data	theory	data	theory	data	theory	data	theory	data	theory	data	theory	data	theory
$\bar{p}/K^-$	$0.27 \pm 0.05$	0.25	$0.32 \pm 0.07$	0.25	$0.31 \pm 0.05$	0.28	$0.10 \pm 0.02$	0.08	$0.07 \pm 0.01$	0.05	$0.04 \pm 0.01$	0.04	$0.02 \pm 0.01$	0.03
$\bar{\Lambda}/K^-$	$0.26 \pm 0.04$	0.21	$0.30 \pm 0.09$	0.20	$0.26 \pm 0.04$	0.20	$0.10 \pm 0.01$	0.08	$0.09 \pm 0.02$	0.05	$0.06 \pm 0.01$	0.04	$0.03 \pm 0.01$	0.03
$\Xi^+/K^-$	$0.04 \pm 0.01$	0.03	$0.04 \pm 0.01$	0.03	$0.03 \pm 0.01$	0.03	$0.018 \pm 0.004$	0.015	$0.018 \pm 0.004$	0.012	$0.009 \pm 0.003$	0.009	$0.006 \pm 0.003$	0.008
$\bar{\Omega}^+/K^-$	—	0.005	—	0.004	$0.005 \pm 0.001$	0.004	$0.004 \pm 0.002$	0.003	—	0.002	—	0.002	—	0.002
$p/K^+$	$0.36 \pm 0.07$	0.33	$0.42 \pm 0.09$	0.36	$0.54 \pm 0.08$	0.44	$1.00 \pm 0.15$	1.03	$1.22 \pm 0.18$	1.32	$2.05 \pm 0.29$	1.79	$1.99 \pm 0.36$	1.93
$\Lambda/K^+$	$0.33 \pm 0.05$	0.25	$0.37 \pm 0.10$	0.26	$0.42 \pm 0.08$	0.28	$0.37 \pm 0.08$	0.51	$0.55 \pm 0.08$	0.60	$0.76 \pm 0.09$	0.67	$0.69 \pm 0.10$	0.72
$\Xi^-/K^+$	$0.04 \pm 0.01$	0.04	$0.04 \pm 0.01$	0.04	$0.04 \pm 0.01$	0.04	$0.049 \pm 0.009$	0.050	$0.050 \pm 0.011$	0.054	$0.057 \pm 0.013$	0.050	$0.055 \pm 0.014$	0.054
$\Omega^-/K^+$	—	0.005	—	0.005	$0.006 \pm 0.001$	0.004	$0.005 \pm 0.001$	0.004	—	0.004	—	0.003	—	0.003

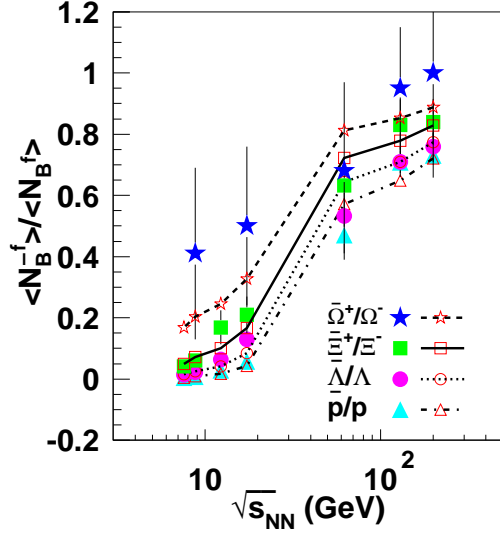


FIG. 1: (Color online) The yield ratios of antibaryons to baryons at mid-rapidity at different energies. The filled symbols with error bars are the experimental results taken from Refs. [24, 26–28, 36]. The corresponding open symbols represent the calculated results where the lines that connect these symbols are just used to guide the eyes.

#### IV. SUMMARY

We study the hadron yield correlations in the combination models. With the basic ideas of the combination mechanism and a few simplifications and/or assumptions based on symmetry and general principles, we show that the hadron yield ratios can be calculated and have a series of regular properties. These ratios are properties of the

TABLE VI: The deduced values of  $d_{\Xi}^{\Lambda}$  and  $d_{\Omega}^p$  at RHIC and SPS energies. Errors shown are total errors. The data are from Refs [24, 26–28, 36].

Energy	$d_{\Xi}^{\Lambda}$	$d_{\Omega}^p$
200 GeV	$0.64 \pm 0.07$	$0.73 \pm 0.23$
130 GeV	$0.59 \pm 0.08$	$0.67 \pm 0.15$
62.4 GeV	$0.34 \pm 0.04$	$0.32 \pm 0.14$
158A GeV	$0.027 \pm 0.009$	$0.029 \pm 0.016$
40A GeV	$0.0013 \pm 0.0005$	$0.0032 \pm 0.0023$

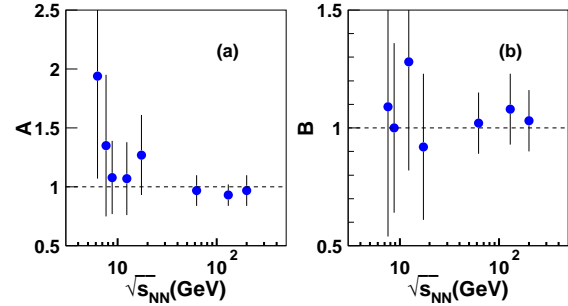


FIG. 2: (Color online) The correlation quantities A (a) and B (b) as the function of the collision energy. The experimental data, filled circles with total statistical and systematic errors, are from Refs. [18, 24–28, 36].

combination mechanism under these assumptions and/or approximations such as the factorization of flavor and momentum dependence of the kernel function, the flavor independence of the momentum distributions, and the approximations made for the net quark contributions. They are

independent of the particular models where usually particular assumptions are made for the kernel functions. A systematic study of these ratios should provide good hints as to whether the combination mechanism is at work. Comparisons with available data are made and predictions for future experiments are given.

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